A Fully-Coupled Approach for Modelling Plastic Deformation and Liquid Lubrication in Metal Forming

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- Theoretical & Numerical Background
- Results & Discussion
- Summary



Theoretical & Numerical Background

- -Equilibrium Equations
- -FE Discretization
- -Viscous Flow within Flow Formulation
- Results & Discussion
- Summary

Equilibrium Equation



In fluid mechanics differential equation of motion is

where $\sigma_{ij} = \sigma_{ij} - p$ and $p = -\sigma_m$

For Newtonian fluid,

$$\tau_{xy} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \qquad \qquad \sigma'_{xx} = 2\mu \frac{\partial u}{\partial x}$$

Equilibrium Equation

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If the space-dependent acceleration term of Navier-Stokes equation is neglected, the equation reduces to

$$\mu \frac{\partial^2 u_i}{\partial x_j^2} - \frac{\partial p}{\partial x_i} + \rho g_i = \rho \frac{du_i}{dt}$$
$$\sigma_{ij} = \sigma'_{ij} + \sigma_m = \sigma'_{ij} - p$$

which gives the dynamic equilibrium equation of solid mechanics that is used in the flow formulation.

$$\frac{\partial \sigma_{ij}}{\partial x_j} + b_i = \rho a_i \qquad b_i = \rho g_i$$
$$a_i = \frac{Du_i}{Dt} = \dot{u}_i + (u_j \cdot \nabla) u_i \cong \dot{u}_i$$

The Flow Formulation

Similarities of viscous flow with flow formulation

• viscous flow formulation

$$\mu \frac{\partial^2 u_i}{\partial x_j^2} - \frac{\partial p}{\partial x_i} + \rho g_i = \rho \frac{du_i}{dt}$$

- identical separation between deviatoric and volumetric terms
- the analogy for 'shear viscosity' and 'bulk viscosity' to ensure the incompressibility constraint of the viscous fluids.

$$\dot{\rho} + \nabla(\rho \mathbf{u}) = 0 \quad \stackrel{\partial \rho / \partial t \cong 0}{\longrightarrow} \quad \nabla \mathbf{u} = 0 \quad \rightarrow \quad \dot{\varepsilon}_{v} = 0$$

$$\sigma_{\rm m} = K \dot{\varepsilon}_{\rm v} \qquad \qquad \sigma_{\rm m} = -p = \mu_{\rm v} \dot{\varepsilon}_{\rm v}$$

 $\sigma_{ij} = \frac{2}{3} \frac{\overline{\sigma}}{\overline{\epsilon}} \dot{\varepsilon}_{ij} \qquad \sigma_{ij} = 2\mu \dot{\varepsilon}_{ij}$

Finite Element Formulation

Weak Form of Equilibrium Equation

$$\begin{split} \partial \Pi &= \int_{V} \sigma_{ij}^{'} \delta \dot{\varepsilon}_{ij}^{'} dV + K \int_{V} \dot{\varepsilon}_{V} \delta \dot{\varepsilon}_{V} dV - \int_{S_{T}} t_{i} \delta u_{i} dS = 0 \\ \sum_{m=1}^{M} \left\{ \int_{V^{m}} \delta \mathbf{v}^{T} K_{\mu} \, \mathbf{v} \, dV^{m} + K^{m} \int_{V^{m}} \delta \mathbf{v}^{T} \mathbf{C}^{T} \mathbf{B} \mathbf{v} \mathbf{C}^{T} \mathbf{B} \, dV^{m} - \int_{S_{T}^{m}} \delta \mathbf{v}^{T} \mathbf{N} \mathbf{T} dS^{m} \right\} = 0 \\ \mathbf{K}_{\mu} &= \mathbf{B}^{T} \mathbf{D}_{\mu} \mathbf{B} \quad \mathbf{C}^{T} = \begin{bmatrix} 1 \ 1 \ 1 \ 0 \ 0 \ 0 \end{bmatrix} \quad \mathbf{B} = \mathbf{L} \, \mathbf{N}^{T} \qquad K^{m} = \mu_{v}^{m} = k_{0} - k_{1} \cdot \frac{\sigma_{ii}}{3} \\ \mathbf{D}_{\mu} &= \begin{bmatrix} 2\mu & 0 & 0 & 0 & 0 \\ 0 & 2\mu & 0 & 0 & 0 \\ 0 & 0 & 2\mu & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{\partial}{\partial y} & 0 \\ 0 & 0 & \frac{\partial}{\partial y} & 0 \\ 0 & \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} & 0 & \frac{\partial}{\partial y} \end{bmatrix} \quad \mathbf{N}^{T} = \begin{bmatrix} N_{1} & 0 & 0 & \dots & N_{8} & 0 & 0 \\ 0 & N_{1} & 0 & \dots & 0 & N_{8} & 0 \\ 0 & 0 & N_{1} & \dots & 0 & 0 & N_{8} \end{bmatrix} \end{split}$$

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Theoretical & Numerical Background

Results & Discussion

- -Bulge Test
- -Lubrication Breakdown
- -Drawing of Sheet with Surface Pockets
- -Deep Drawing

Summary

Bulge Test



Bulge Test



uniform distribution of hydrostatic pressure!

onset of instability (Hill, 1950) $\bar{\varepsilon}_{inst,cal} = \frac{4}{11}(2n+1)$

$$\bar{\varepsilon}_{inst,cal} = 0.39$$

 $\bar{\varepsilon}_{inst,num} = 0.37$

error = 2 %

Lubricant Breakdown





m = 1.0 sticking interfaces

$$\alpha_{lover\ die} = 3.0^{\circ}$$

 $h_{strip} = 1.95\ mm$

Lubricant Breakdown



Shear Stress Normalized by Effective Stress



Drawing of Sheet with Surface Pocket





 $v_{drawing} = 100 \text{ mm/s}$ $\alpha_{lover \, die} = 5.0^{\circ}$ m = 0.3 along the upper toolm = 1.0 sheet - lubricant interface

Blank: AA1050-H111 $\bar{\sigma} = 312\bar{\varepsilon}^{0.08} \,(MPa)$ $\mu = 0.1 Pa \cdot s$ Fluid: lubricant - D 350,0 Hydrostatic pressure [MPa] ••••• E •F 250,0 150,0 50,0 -50.0 210. 40

Displacement [mm]

Deep Drawing



Fluid:

AA1050-H111 $\mu = 10^6 Pa \cdot s$

m = 0.5 sheet – lubricant interface $v_{drawing} = 1 mm/s$





- Theoretical & Numerical Background
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Summary

- -Current State
- -Challenges
- -Future Work

Summary





Future Works









The Flow Formulation

- The weak form of the quasi static stress equilibrium equations used in the flow formulation results from neglecting gravity and dynamic effects.
- Incompressibility is included by transforming the general weak form of the stress equilibrium equations into a weak form with constraints.
- Flow formulation is based on a velocity-pressure approach where metals are considered as fluids of very high viscosity and velocities and hydrostatic pressures are solved simultaneously.

Drawing of Sheet with Surface Pocket



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